

5. Laplace transforms

Laplace transforms are:

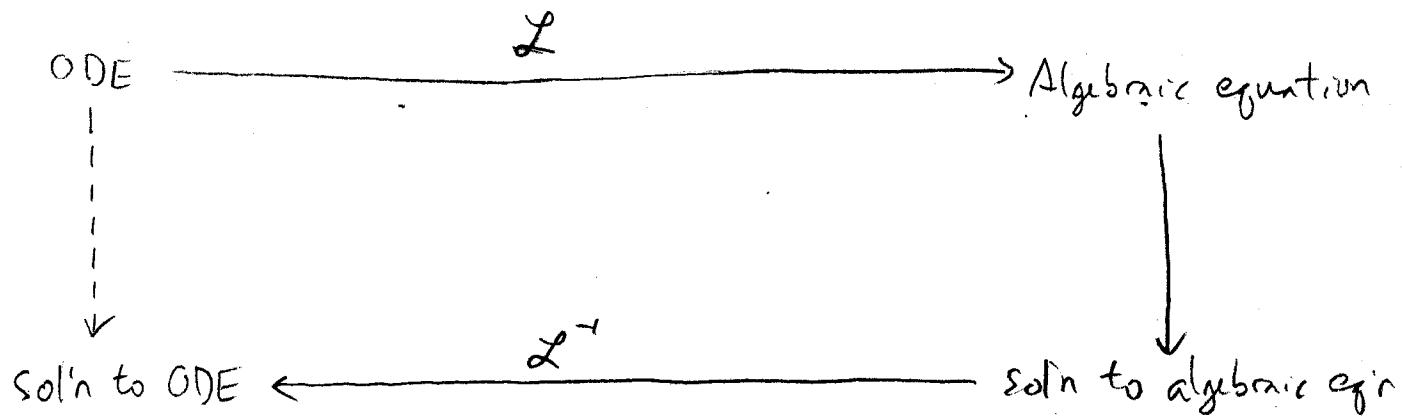
- Used to solve linear ODE's
- Useful when the forcing term is discontinuous,

e.g., step function



(think. Force being turned on/off.)

Big idea:



The Laplace transform is an operator: it inputs a function, and outputs a function.

Def: Suppose $f(t)$ is defined for $0 < t < \infty$. The Laplace transform of f is the function $\mathcal{L}(f)$ where

$$\boxed{\mathcal{L}(f(t))(s) := F(s) = \int_0^\infty f(t) e^{-st} dt, \quad s > 0}$$

Often, we denote $\mathcal{L}(f)$ by F , i.e., $f \xrightarrow{\mathcal{L}} F$

Recall: $\int_0^\infty f(t) dt = \lim_{T \rightarrow \infty} \int_0^T f(t) dt$.

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Example: Compute $\mathcal{L}(f)$, where $f(t) = e^{at}$.

$$\begin{aligned} F(s) &= \int_0^\infty e^{at} e^{-st} dt = \int_0^\infty e^{-(s-a)t} dt = \left. \frac{e^{-(s-a)t}}{-(s-a)} \right|_0^\infty \\ &= \lim_{T \rightarrow \infty} \frac{e^{-(s-a)T}}{-(s-a)} + \frac{1}{s-a} \\ &= \begin{cases} 0 & \text{if } s > a \\ \infty & \text{if } s \leq a \quad (\text{i.e., the limit doesn't exist}) \end{cases} \end{aligned}$$

Thus, $\boxed{\mathcal{L}(e^{at})(s) = \frac{1}{s-a}}$ if $s > a$ (and it isn't defined otherwise).

Remark: Sometimes, the domain of F is restricted.

e.g., $f(t)$ has domain $(-\infty, \infty)$

$F(s)$ has domain (a, ∞)

(Analogy: e^x has domain $(-\infty, \infty)$ but its inverse function $\ln x$ has domain $(0, \infty)$).

Recall: Integration by parts (we'll need it!).

let's rederive it: $(uv)' = u dv + du v$

$$u dv = (uv)' - v du$$

$$\boxed{\int u dv = uv - \int v du}$$

Example: Let $f(t) = t$. Compute $\mathcal{L}(F)$.

$$F(s) = \int_0^\infty t e^{-st} dt. \quad \begin{array}{l} \text{Let } u=t \\ \quad du=dt \end{array} \quad \begin{array}{l} v = -\frac{1}{s} e^{-st} \\ dv = e^{-st} dt \end{array}$$

$$\int \underbrace{t}_{u} \underbrace{e^{-st} dt}_{dv} = \underbrace{-\frac{1}{s} t e^{-st}}_{uv} + \underbrace{\frac{1}{s} \int e^{-st} dt}_{-\int v du} = -\frac{te^{-st}}{s} - \frac{e^{-st}}{s^2}$$

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$$\begin{aligned} \mathcal{L}(f) = F(s) &= \int_0^\infty t e^{-st} dt = \lim_{T \rightarrow \infty} \int_0^T t e^{-st} dt \\ &= \lim_{T \rightarrow \infty} \left(-\frac{t e^{-st}}{s} - \frac{e^{-st}}{s^2} \right) \Big|_0^T = \lim_{T \rightarrow \infty} \left(-\frac{t e^{-st}}{s} - \frac{e^{-st}}{s^2} \right) - \left(0 - \frac{e^0}{s^2} \right) = \boxed{\frac{1}{s^2}} \end{aligned}$$

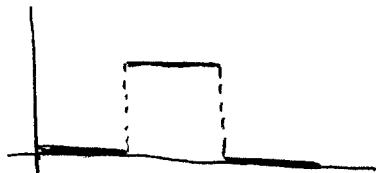
Other common functions: $\mathcal{L}(t^n)(s) = \frac{n!}{s^{n+1}}$

$$\mathcal{L}(\sin bt)(s) = \frac{b}{s^2 + b^2}$$

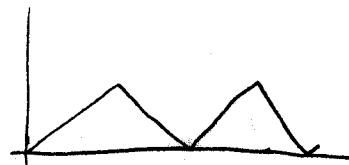
$$\mathcal{L}(\cos bt)(s) = \frac{s}{s^2 + b^2}$$

We can also compute the Laplace transform of piecewise continuous or piecewise differentiable functions.

e.g.,

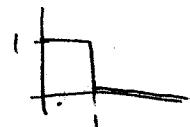


step function
(piecewise continuous)



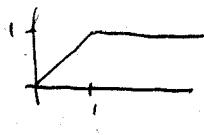
triangle wave
(piecewise differentiable)

Example: Let $f(t) = \begin{cases} 1 & 0 \leq t < 1 \\ 0 & t \geq 1 \end{cases}$



$$\text{Compute } \mathcal{L}(f)(s) = F(s) = \int_0^\infty e^{-st} f(t) dt = \int_0^1 e^{-st} dt = \left[-\frac{e^{-st}}{s} \right]_0^1 = \boxed{-\frac{e^{-s}}{s} + \frac{1}{s}}$$

Example: Let $f(t) = \begin{cases} t & 0 \leq t < 1 \\ 1 & 1 \leq t < \infty \end{cases}$



We must break the integral into 2 parts:

$$\mathcal{L}(f(s)) = F(s) = \int_0^\infty e^{-st} f(t) dt = \underbrace{\int_0^1 t e^{-st} dt}_{I1} + \underbrace{\int_1^\infty e^{-st} dt}_{I2}$$

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$$I_1 = -\frac{e^{-s}}{s} - \left(\frac{e^{-s}}{s^2} - \frac{1}{s^2} \right) \quad I_2 = \lim_{T \rightarrow \infty} \left[\frac{-e^{-st}}{s} \right]_{t=1}^T = \frac{e^{-s}}{s}$$

$$F(s) = \left(-\frac{e^{-s}}{s} - \frac{e^{-s}}{s^2} - \frac{1}{s^2} \right) + \left(\frac{e^{-s}}{s} \right) = \boxed{\frac{1}{s^2} - \frac{e^{-s}}{s^2}}$$

Properties of the Laplace transform:

- \mathcal{L} is linear
- \mathcal{L} "turns derivatives into multiplication."

(i) Linearity: $\mathcal{L}(a f(t) + b g(t))(s) = a \mathcal{L}(f(t))(s) + b \mathcal{L}(g(t))(s)$

i.e., you can break apart sums; pull out constants
(\mathcal{L} is a "linear operator")

(ii) Turns derivatives into multiplication: $\boxed{\mathcal{L}(y'(t))(s) = s Y(s) - y(0)}$

Proof: $\mathcal{L}(y')(s) = \int_0^\infty y'(t) e^{-st} dt = \lim_{T \rightarrow \infty} \left[e^{-st} y(t) + s \int_0^T y(t) e^{-st} dt \right]$
 $= \lim_{T \rightarrow \infty} e^{-sT} y(T) \Big|_0^T + s \mathcal{L}(y)(s)$
 $= \underbrace{\lim_{T \rightarrow \infty} e^{-sT} y(T)}_{=0, \text{ as long as } |y(t)| \leq C e^{at}} - y(0) + s \mathcal{L}(y)(s) = s \mathcal{L}(y)(s) - y(0) = s Y(s) - y(0).$ ✓

$|y(t)| \leq C e^{at}$. Henceforth, we will make this blanket assumption.

Similarly, $\mathcal{L}(y'')(s) = s^2 Y(s) - s y(0) - y'(0)$

$$\mathcal{L}(y''')(s) = s^3 Y(s) - s^2 y(0) - s y'(0) - y''(0)$$

$$\mathcal{L}(y^{(4)})(s) = s^4 Y(s) - s^3 y(0) - s^2 y'(0) - s y''(0) - y'''(0).$$

⋮
etc

$$\begin{aligned}
 \text{Proof: } \mathcal{L}(y'') &= s\mathcal{L}(y') - y(0) \\
 &= s(s\mathcal{L}(y) - y(0)) - y'(0) \\
 &= s^2 Y - s y(0) - y'(0) \quad \checkmark
 \end{aligned}$$

The proofs of the formulas for higher derivatives is handled similarly. □

More Laplace transform properties

$$(i) \mathcal{L}(e^{ct} f(t))(s) = F(s-c)$$

$$(ii) \mathcal{L}(t f(t))(s) = -F'(s)$$

$$(iii) \mathcal{L}(t^n f(t))(s) = (-1)^n F^{(n)}(s)$$

Applications/examples of this:

Example 1: $f(t) = e^{2t} \cos 3t$

Recall $\mathcal{L}(\cos 3t) = \frac{s}{s^2 + 9}$.

Using (i), $\mathcal{L}(e^{2t} \cos 3t) = \boxed{\frac{s-2}{(s-2)^2 + 9}} = F(s)$.

Example 2: $f(t) = t^2 e^{3t}$. let $g(t) = e^{3t}$

Use (ii): Recall: $\mathcal{L}(t^2) = \frac{2}{s^3} \Rightarrow$ by (i), $\mathcal{L}(t^2 e^{3t}) = \boxed{\frac{2}{(s-3)^3}}$

OR Use (iii): $\mathcal{L}(t^2 e^{3t}) = (-1)^2 F''(s) = \frac{d^2}{ds^2} \left(\frac{1}{s-3} \right) = \frac{2}{(s-3)^3}$.

Using the Laplace transform to solve ODEs

Example: Consider the IVP: $y'' - y = e^{2t}$, $y(0) = 0$, $y'(0) = 1$.

Apply \mathcal{L} : $\mathcal{L}(y'') - \mathcal{L}(y) = \mathcal{L}(e^{2t})$.

$$[s^2 Y - s \overset{(0)}{y} - y'(0)] - Y = \frac{1}{s-2}$$

$$s^2 Y - 1 - Y = \frac{1}{s-2} \Rightarrow (s^2 - 1)Y = \frac{1}{s-2} + \frac{s-2}{s-2} = \frac{s-1}{s-2}$$

$$\Rightarrow \boxed{Y(s) = \frac{1}{(s+1)(s-2)}}$$

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* The solution to the IVP
is the function whose
Laplace transform is

$$Y(s) = \frac{1}{(s+1)(s-2)}$$

To do this, write $\frac{1}{(s+1)(s-2)} = \frac{A}{(s+1)} + \frac{B}{(s-2)}$

Use partial fraction decomposition:

$$\frac{A}{(s+1)} \frac{(s-2)}{(s-2)} + \frac{B}{(s-2)} \frac{(s+1)}{(s+1)} = \frac{1}{(s+1)(s-2)} \Rightarrow \begin{cases} A=0 \\ B=1 \end{cases} \Rightarrow (A+B)s + (B-2A) = 1$$

$$\Rightarrow \begin{cases} A+B=0 \\ B-2A=1 \end{cases} \Rightarrow A=-B \Rightarrow 3B=1 \Rightarrow \begin{cases} B=\frac{1}{3} \\ A=-\frac{1}{3} \end{cases}$$

$$\text{So, } \frac{1}{(s+1)(s-2)} = \frac{-\frac{1}{3}}{s+1} + \frac{\frac{1}{3}}{s-2}$$

$$\begin{aligned} \Rightarrow \mathcal{L}^{-1}\left(\frac{1}{(s+1)(s-2)}\right) &= \mathcal{L}^{-1}\left(\frac{-\frac{1}{3}}{s+1}\right) + \mathcal{L}^{-1}\left(\frac{\frac{1}{3}}{s-2}\right) \\ &= -\frac{1}{3} \mathcal{L}^{-1}\left(\frac{1}{s+1}\right) + \frac{1}{3} \mathcal{L}^{-1}\left(\frac{1}{s-2}\right) = \boxed{-\frac{1}{3}e^{-t} + \frac{1}{3}e^{2t}} \end{aligned}$$

Example (when partial fractions fails).

Compute $\mathcal{L}^{-1}\left(\frac{1}{s^2+4s+13}\right)$. Problem: $s^2+4s+13$ doesn't factor.

Instead, put it in the form $\frac{1}{(s-a)^2+b^2}$; because $\mathcal{L}(e^{at} \sin bt) = \frac{b}{(s-a)^2+b^2}$.

"Complete the square:"

$$\frac{1}{(s^2+4s+4)+9} = \frac{1}{(s+2)^2+3^2} = \frac{1}{3} \frac{3}{(s+2)^2+3^2} \xrightarrow{\mathcal{L}^{-1}} \boxed{\frac{1}{3} e^{-2t} \sin 3t}$$

Example: Solve the IVP: $y'' - 2y' - 3y = 0$, $y(0) = 1$, $y'(0) = 0$.

Old method: $y(t) = e^{rt}$, ... $e^{rt}(r^2 - 2r - 3) = 0$

$$\Rightarrow (r-3)(r+1) = 0 \Rightarrow y(t) = C_1 e^{3t} + C_2 e^{-t}$$

use IC's: $y(0) = C_1 + C_2 = 1$ and $y'(t) = 3C_1 e^{3t} - C_2 e^{-t}$
 $y'(0) = 3C_1 - C_2 = 0$

$$\Rightarrow \begin{cases} C_1 + C_2 = 1 \\ 3C_1 - C_2 = 0 \end{cases} \Rightarrow C_1 = \frac{1}{4}, \quad C_2 = \frac{3}{4} \Rightarrow y(t) = \frac{1}{4} e^{3t} + \frac{3}{4} e^{-t}$$

New method: $y'' - 2y' - 3y = 0$, $y(0) = 1$, $y'(0) = 0$.

$$\mathcal{L}(y'') - 2\mathcal{L}(y') - 3\mathcal{L}(y) = 0$$

$$[s^2 Y - s y(0) - y'(0)] - 2[sY - y(0)] - 3Y = 0.$$

$$[s^2 Y - s - 0] - 2[sY - 1] - 3Y = 0 \Rightarrow (s^2 - 2s - 3)Y = s - 2$$

$$Y = \frac{s-2}{s^2 - 2s - 3} = \frac{A}{s-3} + \frac{B}{s+1} = \frac{s-2}{s^2 - 2s - 3}$$

$$\frac{A}{s-3} \frac{(s+1)}{(s+1)} + \frac{B}{s+1} \frac{(s-3)}{(s-3)} = \frac{(A+B)s + (A-3B)}{(s+1)(s-3)} = \frac{s-2}{(s+1)(s-3)}$$

$$\begin{cases} A+B=1 \\ A-3B=-2 \end{cases} \Rightarrow \begin{cases} A=\frac{1}{4} \\ B=\frac{3}{4} \end{cases} \Rightarrow Y(s) = \frac{1/4}{s-3} + \frac{3/4}{s+1}$$

$$y(t) = \mathcal{L}^{-1}\left(\frac{1/4}{s-3}\right) + \mathcal{L}^{-1}\left(\frac{3/4}{s+1}\right) = \frac{1}{4} \mathcal{L}^{-1}\left(\frac{1}{s-3}\right) + \frac{3}{4} \mathcal{L}^{-1}\left(\frac{1}{s+1}\right)$$

$$\Rightarrow \boxed{y(t) = \frac{1}{4} e^{3t} + \frac{3}{4} e^{-t}}$$

Summary / analysis: Consider $ay'' + by' + cy = f(t)$, $y(0) = y_0$, $y'(0) = y_1$.

$$\begin{aligned} \mathcal{L}(ay'' + by' + cy) &= a\mathcal{L}(y'') + b\mathcal{L}(y') + c\mathcal{L}(y) \\ &= a(s^2 Y - s y(0) - y'(0)) + b(sY - y(0)) + cY \\ &= (as^2 + bs + c)Y - y_0(as + b) - ay_1 = F(s). \end{aligned}$$

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$$\text{Thus, } Y(s) = \underbrace{\frac{F(s)}{as^2 + bs + c}}_{\text{"state-free sol'n"}} + \underbrace{\frac{(as+b)y_0 + ay_1}{as^2 + bs + c}}_{\text{"input-free sol'n"}}$$

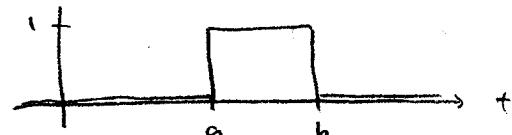
i.e., $Y(s) = Y_s(s) + Y_i(s)$, where

$Y_s(s)$ doesn't depend on the initial conditions, $y(0)=y_0$, $y'(0)=y_1$.

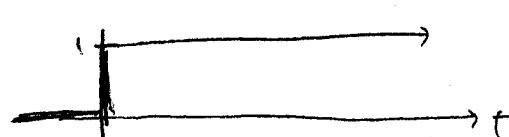
$Y_i(s)$ doesn't depend on the forcing term $f(t)$.

Discontinuous forcing terms:

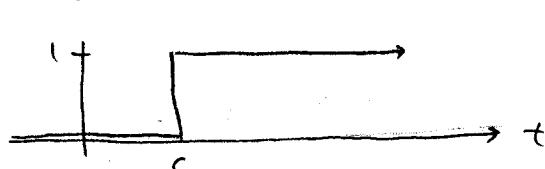
- Interval function: $H_{ab}(t) = \begin{cases} 0 & t < a \\ 1 & a \leq t < b \\ 0 & b \leq t < \infty \end{cases}$



- Heaviside function: $H(t) = \begin{cases} 0 & t < 0 \\ 1 & t \geq 0 \end{cases}$

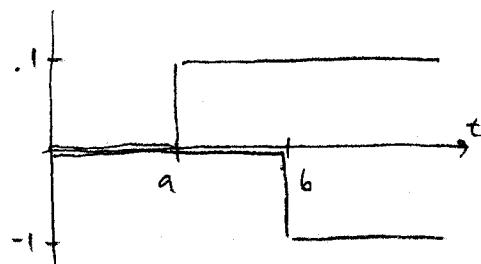


- Shifted Heaviside function: $H_c(t) = \begin{cases} 0 & t < c \\ 1 & t \geq c \end{cases}$



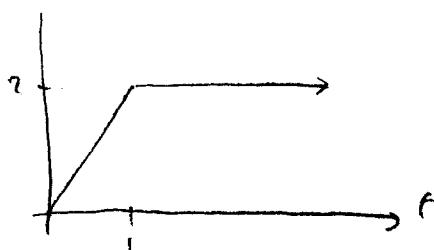
Remarks: • $H_c(t) = H(t-c)$

• $H_{ab}(t) = H_a(t) - H_b(t) = H(t-a) - H(t-b)$



* Many piecewise continuous functions can be written using Heaviside functions

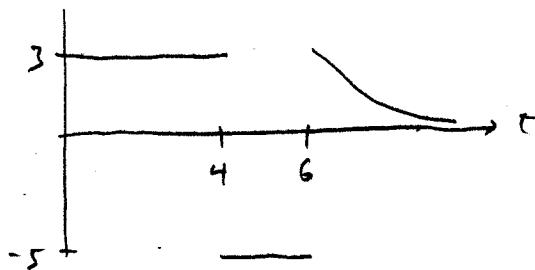
Example 1: $f(t) = \begin{cases} 2t & 0 \leq t < 1 \\ 2 & t \geq 1 \end{cases}$ $f(t) = 2t H_{01}(t) + 2 H_1(t)$



$$= 2t [H(t) - H(t-1)] + 2 H(t-1)$$

$$= \boxed{2t H(t) - 2(t-1) H(t-1)}$$

Example 2: $f(t) = \begin{cases} 3 & 0 \leq t < 4 \\ -5 & 4 \leq t < 6 \\ e^{7-t} & 6 \leq t < \infty \end{cases}$



$$\begin{aligned} f(t) &= 3H_{04}(t) - 5H_{46}(t) + e^{7-t}H_6(t) \\ &= 3[H(t) - H(t-4)] - 5[H(t-4) - H(t-6)] + e^{7-t}[H(t-6)] \\ &= 3H(t) - 8H(t-4) + 5H(t-6) + e^{7-t}H(t-6) \end{aligned}$$

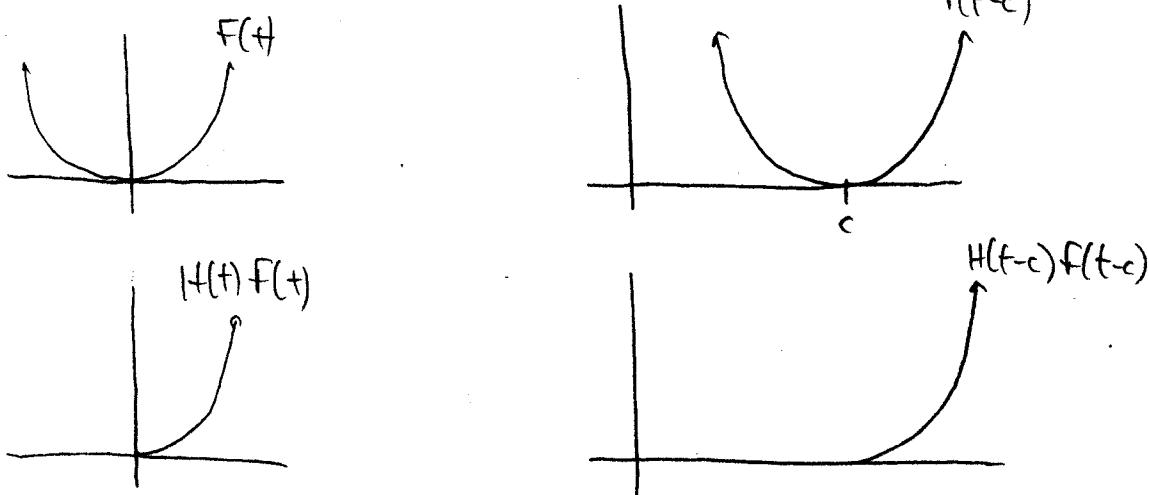
Motivation: Writing piecewise continuous functions in this form is convenient for taking their Laplace transform.

Recall: • $\mathcal{L}(H_{ab}(t)) = \int_a^b e^{-st} dt = \frac{e^{-as} - e^{-bs}}{s}$

• $\mathcal{L}(e^{ct} f(t)) = F(s-c)$

Fact: $\boxed{\mathcal{L}(H(t-c)f(t-c)) = e^{-cs}F(s)}$

How to interpret this:



Thus, $H(t-c)f(t-c)$ truncates $f(t)$ at $t=0$, then shifts it by c . Since $f(t)$ is defined for only $t \geq 0$, we don't want $f(t-c)$ to be defined for $t < c$, so we use $H(t-c)f(t-c)$.

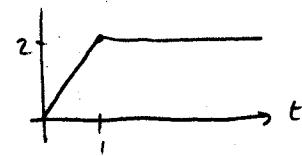
[10]

Example: Compute the Laplace transform of $f(t) = \begin{cases} 2t & 0 \leq t < 1 \\ 2 & t \geq 1 \end{cases}$.

Recall: $f(t) = 2tH(t) - 2(t-1)H(t-1)$.

$$\mathcal{L}\{f(t)\} = 2\mathcal{L}\{tH(t)\} - 2\mathcal{L}\{(t-1)H(t-1)\}$$

$$= \boxed{\frac{2}{s^2} - \frac{2}{s^2}e^{-s}}$$



More practice using $\mathcal{L}\{H(t-c)f(t-c)\} = e^{-cs}F(s)$:

- $\mathcal{L}\{(t-3)^2H(t-3)\} = e^{-3s}F(s)$

$$f(t-3) = (t-3)^2, \text{ so } f(t) = f((t+3)-3) = ((t+3)-3)^2 = t^2$$

$$F(s) = \frac{1}{s^2} \Rightarrow \mathcal{L}\{(t-3)^2H(t-3)\} = e^{-3s}F(s) = \boxed{\frac{e^{-3s}}{s^2}}$$

- $\mathcal{L}\{t^2H(t-3)\} = e^{-3s}F(s)$

$$f(t-3) = t^2, \text{ so } f(t) = f((t+3)-3) = (t+3)^2 = t^2 + 6t + 9$$

$$F(s) = \frac{2}{s^3} + \frac{6}{s^2} + \frac{9}{s} \Rightarrow \mathcal{L}\{t^2H(t-3)\} = e^{-3s}F(s) = \boxed{e^{-3s}\left(\frac{2}{s^3} + \frac{6}{s^2} + \frac{9}{s}\right)}$$

- $\mathcal{L}\{e^{t-1}H(t-1)\} = e^{-s}F(s)$

$$f(t-1) = e^{t-1}, \text{ so } F(t) = F((t+1)-1) = e^{(t+1)-1} = e^t$$

$$F(s) = \frac{1}{s-1} \Rightarrow \mathcal{L}\{e^{t-1}H(t-1)\} = \boxed{e^{-s} \frac{1}{s-1}}$$

- $\mathcal{L}\{e^{7-t}H(t-6)\} = e^{-6s}F(s)$

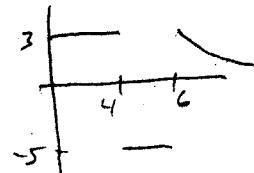
$$f(t-6) = e^{7-t}, \text{ so } F(t) = F((t+6)-6) = e^{7-(t+6)} = e^{1-t} = e \cdot e^{-t}$$

$$F(s) = e \cdot \frac{1}{s+1} \Rightarrow \mathcal{L}\{e^{7-t}H(t-6)\} = e^{-6s}F(s) = e^{-6s} \cdot e \frac{1}{s+1} = \boxed{\frac{e^{1-6s}}{s+1}}$$

Trick: Given $f(t-c)$, plug in $t+c$ for t to find $F(s)$,

e.g., $f((t+c)-c) = f(t)$.

Example: Find $F(s)$, where $f(t) = \begin{cases} 3 & 0 \leq t < 4 \\ -5 & 4 \leq t < 6 \\ e^{7-t} & 6 \leq t < \infty \end{cases}$



Recall: $f(t) = 3H(t) - 8H(t-4) + 5H(t-6) + e^{7-t}H(t-6)$

$$F(s) = \frac{3}{s} - \frac{8}{s} e^{-4s} + \frac{5}{s} e^{-6s} + \frac{1}{s+1} e^{1-6s} = \frac{3-8e^{-4s}+5e^{-6s}}{s} + \frac{e^{1-6s}}{s+1}$$

Example: Solve the IVP: $y'' + y = f(t)$, $y(0) = 0$, $y'(0) = 1$,

where $f(t) = \begin{cases} 2t & 0 \leq t \leq 1 \\ 2 & t > 1 \end{cases}$

Recall: $f(t) = 2tH(t) - 2(t-1)H(t-1)$.

$$F(s) = \frac{2}{s^2} - \frac{2}{s^2} e^{-s}$$

Take \mathcal{L} of both sides of the ODE:

$$[s^2 Y - s y(0) - y'(0)] + Y = \frac{2-2e^{-s}}{s^2}$$

$$s^2 Y - 1 + Y = \frac{2-2e^{-s}}{s^2}$$

$$(s^2 + 1)Y = \frac{2-2e^{-s}}{s^2} + 1 \Rightarrow Y(s) = \frac{2-2e^{-s}}{s^2(s^2+1)} + \frac{1}{s^2+1}$$

Partial fractions: $\frac{1}{s^2(s^2+1)} = \frac{1}{s^2} - \frac{1}{s^2+1}$

$$\begin{aligned} Y(s) &= \frac{2}{s^2} - \frac{2}{s^2+1} - \frac{2e^{-s}}{s^2} + \frac{2e^{-s}}{s^2+1} + \frac{1}{s^2+1} \\ &= \frac{2}{s^2} - \frac{2e^{-s}}{s^2} - \frac{1}{s^2+1} + \frac{2e^{-s}}{s^2+1} \end{aligned}$$

$$\begin{aligned} y(t) &= 2\mathcal{L}^{-1}\left(\frac{1}{s^2}\right) - 2\mathcal{L}^{-1}\left(\frac{e^{-s}}{s^2}\right) - \mathcal{L}^{-1}\left(\frac{1}{s^2+1}\right) + 2\mathcal{L}^{-1}\left(\frac{e^{-s}}{s^2+1}\right) \\ &= 2t - 2(t-1)H(t-1) - \sin t + 2\sin(t-1)H(t-1) \\ &= [2t - \sin t] + [2\sin(t-1) - 2(t-1)]H(t-1) \\ &= \begin{cases} 2t - \sin t & 0 \leq t < 1 \\ 2 + 2\sin(t-1) - \sin t & t \geq 1 \end{cases} \end{aligned}$$

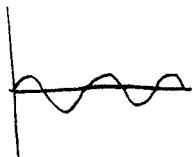
* This is the unique sol'n to the IVP $y'' + y = f(t)$, $y(0) = 0$, $y'(0) = 1$.

②

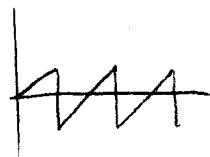
Periodic forcing terms:

- Suppose $f(t)$ is periodic. We want to compute $F(s) = \mathcal{L}(f(t))(s)$.

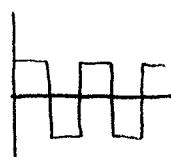
e.g.,



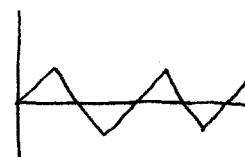
sine wave



sawtooth wave



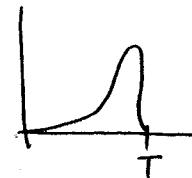
square wave



triangle wave.

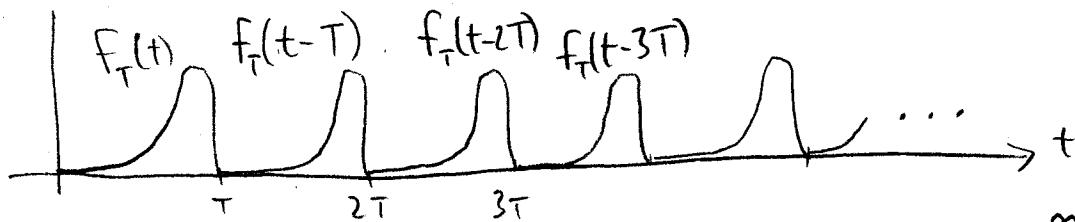
Approach: Consider the "window" $f_T(t)$, defined on

$0 \leq t < T$, and extended to be periodic



Then, the function $f(t) = \begin{cases} f_T(t) & 0 \leq t < T \\ f_T(t-kT) & kT \leq t < (k+1)T \end{cases}$

is periodic:



Remark: $f_T(t-kT) = f_T(t-kT)H(t-kT)$, so $f(t) = \sum_{k=0}^{\infty} f_T(t-kT)H(t-kT)$.

key point: If $|x| < 1$, then $1+x+x^2+x^3+\dots = \boxed{\sum_{n=0}^{\infty} x^n = \frac{1}{1-x}}$

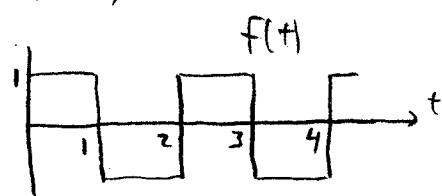
$$\text{Apply this to } F(s) = \sum_{k=0}^{\infty} \mathcal{L}\{f_T(t-kT)H(t-kT)\} = \sum_{k=0}^{\infty} e^{-kTs} F_T(s)$$

$$= F_T(s) \sum_{k=0}^{\infty} e^{-kTs} = F_T(s) \sum_{k=0}^{\infty} (e^{-sT})^k = \boxed{F_T(s) \frac{1}{1-e^{-sT}}}$$

Example: Solve the IVP $y'' + y = f(t)$, $y(0) = 0$, $y'(0) = 0$,

where $f(t)$ is the following square wave,

of period $T = 2$.



First, compute $F(s)$:

$$\begin{aligned} F_T(t) &= H_{01}(t) - H_{12}(t) = [H(t) - H(t-1)] - [H(t-1) - H(t-2)] \\ &= H(t) - 2H(t-1) + H(t-2) \end{aligned}$$

$$\begin{aligned} F_T(s) &= \mathcal{L}(H(t)) - 2\mathcal{L}(H(t-1)) + \mathcal{L}(H(t-2)) \\ &= \frac{1}{s} - \frac{2}{s}e^{-s} + \frac{1}{s}e^{-2s} = \boxed{\frac{(1-e^{-s})^2}{s}} \end{aligned}$$

$$\text{Now, } F(s) = F_T(s) \frac{1}{1-e^{-2s}} = \frac{F_T(s)}{(1-e^{-s})(1+e^{-s})} = \frac{(1-e^{-s})(1-e^{-s})}{s(1-e^{-s})(1+e^{-s})} = \boxed{\frac{(1-e^{-s})}{s(1+e^{-s})}}$$

Back to the IVP: $\mathcal{L}(y'') + \mathcal{L}(y) = F(s)$

$$\begin{aligned} [s^2Y - s y(0) - y'(0)] + Y &= \frac{1-e^{-s}}{s(1+e^{-s})} \\ (s^2+1)Y &= \frac{1-e^{-s}}{s(1+e^{-s})} \Rightarrow \boxed{Y(s) = \frac{1-e^{-s}}{s(s^2+1)(1+e^{-s})}} \end{aligned}$$

Simplify this: $Y(s) = \underbrace{\frac{1}{s(s^2+1)}}_{\frac{1}{s} - \frac{s}{s^2+1}} \cdot \underbrace{\frac{1-e^{-s}}{1+e^{-s}}}_{\text{need to further simplify.}}$

$$\frac{1-e^{-s}}{1+e^{-s}} = -\frac{(1+e^{-s})}{1+e^{-s}} + \frac{2}{1+e^{-s}} = 2\left(\frac{1}{1-(-e^{-s})}\right) = 2 \sum_{n=0}^{\infty} (-1)^n e^{-ns}$$

$$\text{Using } \frac{1}{1-x} = \sum_{n=0}^{\infty} x^n, \text{ we have } \frac{1}{1+x} = \frac{1}{1-(-x)} = \sum_{n=0}^{\infty} (-x)^n = \sum_{n=1}^{\infty} (-1)^n x^n.$$

$$\begin{aligned} \text{Thus, } Y(s) &= \left(\frac{1}{s} - \frac{s}{s^2+1}\right) \left(-1 + 2 \sum_{n=0}^{\infty} (-1)^n e^{-ns}\right) \\ &= \left(\frac{1}{s} - \frac{s}{s^2+1}\right) \left(+1 + 2 \sum_{n=1}^{\infty} (-1)^n e^{-ns}\right) = G(s) + 2 \sum_{n=1}^{\infty} (-1)^n e^{-ns} G(s). \end{aligned}$$

Call this $G(s)$. Note that $\boxed{g(t) = 1 - \cos t}$.

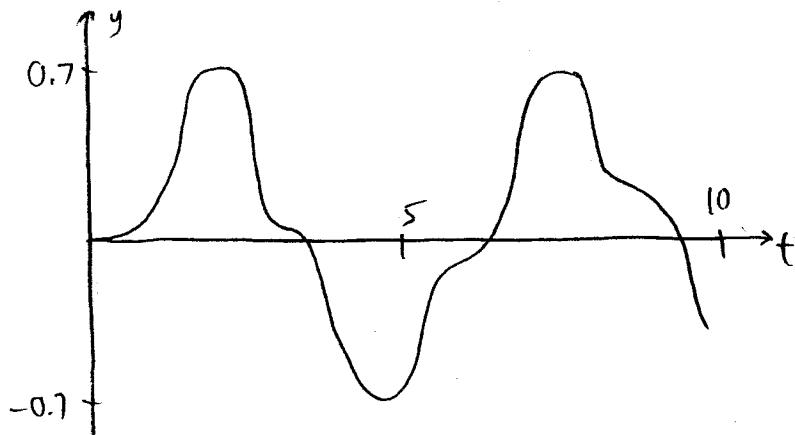
$$\text{Apply } \mathcal{L}\{g(t-n)H(t-n)\} = e^{-ns} G(s).$$

$$Y(s) = G(s) + 2 \sum_{n=1}^{\infty} (-1)^n e^{-ns} G(s) \implies y(t) = g(t) + 2 \sum_{n=1}^{\infty} (-1)^n g(t-n) H(t-n)$$

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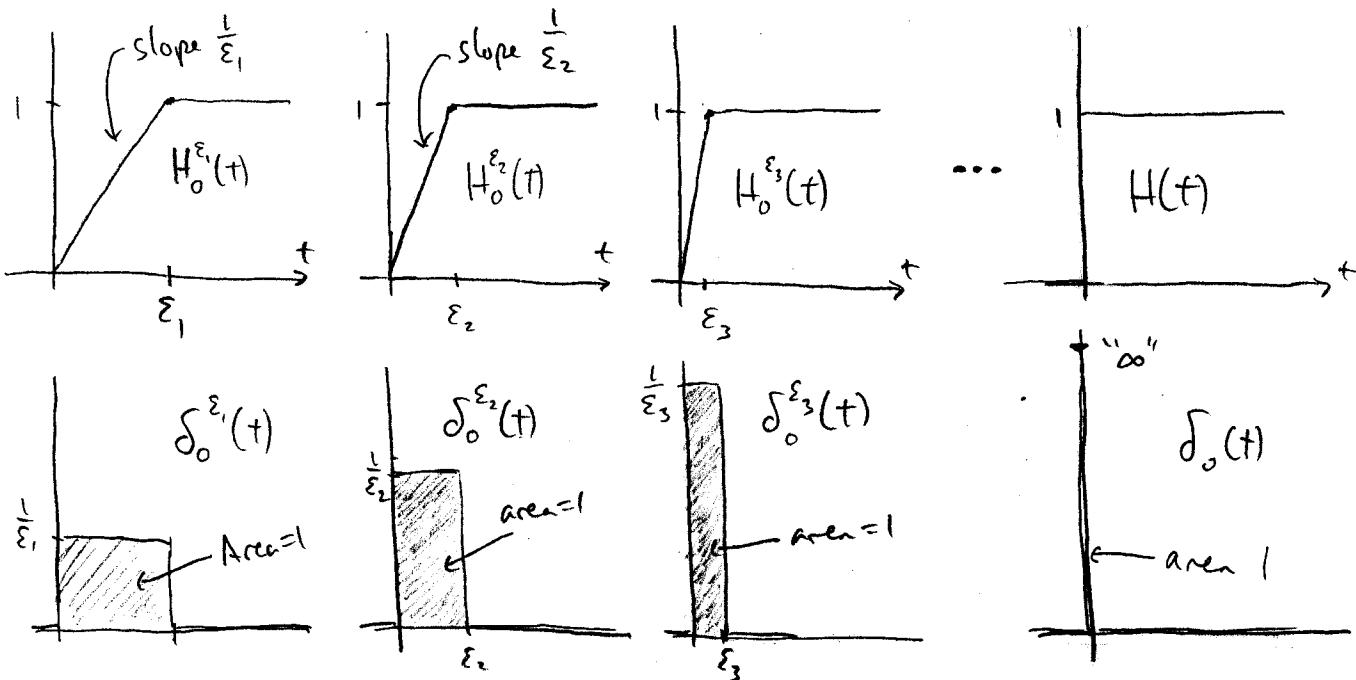
$$y(t) = (1 - \cos t) H(t) + 2 \sum_{n=1}^{\infty} (-1)^n [1 - \cos(t-n)] H(t-n)$$

This is a superposition of infinitely many waves:



Question: What is the "derivative" of the Heavyside function?

Technically, it's not defined, but what "should" it be?



Def: The delta function is $\delta_p(t) = \lim_{\epsilon \rightarrow 0} \delta_p^\epsilon(t)$; $\delta_p(t) = \begin{cases} 0 & t \neq p \\ \infty & t = p \end{cases}$

Technically, it's not really a function, but it's useful!

It has infinite height, infinitesimal width, and integral 1.

Properties of the Delta function:

- $\int_{-\infty}^{\infty} \delta_p(t) dt = 1 \Rightarrow \int_{-\infty}^{\infty} \delta_p(t) f(t) dt = f(p)$

- $\mathcal{L}(\delta_0) = 1 \Rightarrow \mathcal{L}(\delta_p) = e^{-sp}$

So now we can take the inverse Laplace transform of a constant or an exponential.

- * The Delta function models a unit impulse force (finite force over an infinitesimal time interval (or realistically, a very small interval)).
e.g., Exerting a force by hitting something with a hammer.

Example: Solve the IVP $y'' + 2y' + 2y = \delta_0(t)$, $y(0) = 0$, $y'(0) = 0$.

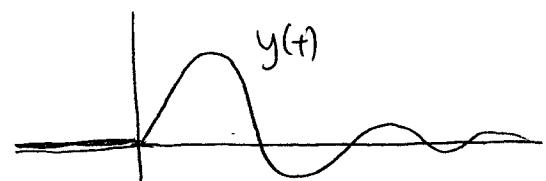
First, take the Laplace transform: $\mathcal{L}(y'' + 2y' + 2y) = \mathcal{L}(\delta_0(t))$
 $(s^2 + 2s + 2)Y = 1$

$$\Rightarrow Y(s) = \frac{1}{s^2 + 2s + 2} = \frac{1}{(s+1)^2 + 1} \Rightarrow \boxed{y(t) = e^{-t} \sin t}$$

Remark: $y(0) = 0$ but $y'(0) = 1$. (This doesn't match the initial conditions!)

This is okay, because $y(t)$ is only defined for $t > 0$, so what we really have

$$\text{is } H(t) y(t) = \begin{cases} 0 & t \leq 0 \\ e^{-t} \sin t & t > 0 \end{cases}$$



This isn't even differentiable at $t=0$, so technically, the derivative isn't defined. But $\lim_{t \rightarrow 0^-} \frac{d}{dt} (H(t) y(t)) = 0$, and that's "good enough!"