# Lecture 4.5: Generalized Fourier series

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## Last time

#### Definition

A Sturm-Liouville equation is a 2nd order ODE of the following form:

$$-(p(x)y')' + q(x)y = \lambda w(x)y, \qquad \text{where } p(x), \ q(x), \ w(x) > 0.$$

We are usually interested in solutions y(x) on a bounded interval [a, b], under some homogeneous BCs:

$$\begin{aligned} \alpha_1 y(a) + \alpha_2 y'(a) &= 0 \qquad \alpha_1^2 + \alpha_2^2 > 0 \\ \beta_1 y(b) + \beta_2 y'(b) &= 0 \qquad \beta_1^2 + \beta_2^2 > 0. \end{aligned}$$

Together, this BVP is called a Sturm-Liouville (SL) problem.

## Main theorem

Given a Sturm-Liouville problem:

- (a) The eigenvalues are real and can be ordered so  $\lambda_1 < \lambda_2 < \lambda_3 < \cdots \rightarrow \infty$ .
- (b) Each eigenvalue  $\lambda_i$  has a unique (up to scalars) eigenfunction  $y_i(x)$ .
- (c) W.r.t. the inner product  $\langle f, g \rangle := \int_a^b f(x)\overline{g(x)}w(x) dx$ , the eigenfunctions form an orthonormal basis on the subspace of functions  $C_{\alpha,\beta}^{\infty}[a, b]$  that satisfy the BCs.

## What this means

#### Main theorem

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#### Definition

If  $f \in C^{\infty}_{\alpha,\beta}[a, b]$ , then f can be written uniquely as a linear combination of the eigenfunctions. That is,

$$f(x) = \sum_{n=1}^{\infty} c_n y_n(x), \quad \text{where } c_n = \frac{\langle f, y_n \rangle}{\langle y_n, y_n \rangle} = \frac{\int_a^b f(x) \overline{y_n(x)} w(x) \, dx}{\int_a^b ||y_n(x)||^2 w(x) \, dx}$$

This is called a generalized Fourier series with respect to the orthogonal basis  $\{y_n(x)\}$  and weighting function w(x).

## Example 1 (Dirichlet BCs)

$$-y^{\prime\prime}=\lambda y, \hspace{1em} y(0)=0, \hspace{1em} y(\pi)=0 \hspace{1em}$$
 is an SL problem with:

• Eigenvalues: 
$$\lambda_n = n^2$$
,  $n = 1, 2, 3, \ldots$ 

• Eigenfunctions: 
$$y_n(x) = \sin(nx)$$
.

The orthogonality of the eigenvectors means that

$$\langle y_m, y_n \rangle := \int_0^{\pi} y_m(x) y_n(x) w(x) \, dx = \int_0^{\pi} \sin(mx) \, \sin(nx) \, dx = \begin{cases} 0 & \text{if } m \neq n \\ \pi/2 & \text{if } m = n. \end{cases}$$

Note that this means that  $||y_n|| := \langle y_n, y_n \rangle^{1/2} = \sqrt{\pi/2}$ .

Fourier series: any function f(x), continuous on  $[0, \pi]$  satisfying f(0) = 0,  $f(\pi) = 0$  can be written *uniquely* as

$$f(x) = \sum_{n=1}^{\infty} b_n \sin nx$$

where

$$b_n = \frac{\langle f, \sin nx \rangle}{\langle \sin nx, \sin nx \rangle} = \frac{\int_0^\pi f(x) \sin nx \, dx}{||\sin nx||^2} = \frac{2}{\pi} \int_0^\pi f(x) \sin nx \, dx.$$

## Example 2 (Neumann BCs)

$$-y^{\prime\prime} = \lambda y$$
,  $y^{\prime}(0) = 0$ ,  $y^{\prime}(\pi) = 0$  is an SL problem with:

• Eigenvalues: 
$$\lambda_n = n^2$$
,  $n = 0, 1, 2, 3, ...$ 

• Eigenfunctions:  $y_n(x) = \cos(nx)$ .

The orthogonality of the eigenvectors means that

$$\langle y_m, y_n \rangle := \int_0^{\pi} y_m(x) y_n(x) w(x) \, dx = \int_0^{\pi} \cos(mx) \, \cos(nx) \, dx = \begin{cases} 0 & \text{if } m \neq n \\ \pi/2 & \text{if } m = n > 0. \end{cases}$$

Note that this means that  $||y_n|| := \langle y_n, y_n \rangle^{1/2} = \begin{cases} \sqrt{\pi/2} & n > 0\\ \sqrt{\pi} & n = 0. \end{cases}$ 

Fourier series: any function f(x), continuous on  $[0, \pi]$  satisfying f'(0) = 0,  $f'(\pi) = 0$  can be written *uniquely* as

$$f(x) = \sum_{n=0}^{\infty} a_n \cos nx$$

where

$$a_n = \frac{\langle f, \cos nx \rangle}{\langle \cos nx, \cos nx \rangle} = \frac{\int_0^\pi f(x) \cos nx \, dx}{||\cos nx||^2} = \frac{2}{\pi} \int_0^\pi f(x) \cos nx \, dx.$$

The same formula holds for  $a_0$  if you let the n = 0 (constant) term be  $\frac{a_0}{2}$  rather than  $a_0$ .

## Example 3 (Mixed BCs)

$$-y'' = \lambda y$$
,  $y(0) = 0$ ,  $y'(\pi) = 0$  is an SL problem with:

• Eigenvalues: 
$$\lambda_n = \left(n + \frac{1}{2}\right)^2$$
,  $n = 0, 1, 2, ...$ 

• Eigenfunctions: 
$$y_n(x) = \sin(n + \frac{1}{2})x$$
.

The orthogonality of the eigenvectors means that

$$\langle y_m, y_n \rangle := \int_0^{\pi} \sin(m + \frac{1}{2}) x \sin(n + \frac{1}{2}) x w(x) dx = \begin{cases} 0 & \text{if } m \neq n \\ \pi/2 & \text{if } m = n. \end{cases}$$

Note that this means that  $||y_n|| := \langle y_n, y_n \rangle^{1/2} = \sqrt{\pi/2}$ .

(Generalized?) Fourier series: any function f(x), continuous on  $[0, \pi]$  satisfying f(0) = 0,  $f'(\pi) = 0$  can be written *uniquely* as

$$f(x) = \sum_{n=1}^{\infty} b_n \sin\left(n + \frac{1}{2}\right) x$$

where

$$b_n = \frac{\langle f, \sin(n+\frac{1}{2})x \rangle}{\langle \sin(n+\frac{1}{2})x, \sin(n+\frac{1}{2})x \rangle} = \frac{\int_0^{\pi} f(x)\sin(n+\frac{1}{2})x \, dx}{||\sin(n+\frac{1}{2})x||^2} = \frac{2}{\pi} \int_0^{\pi} f(x)\sin(n+\frac{1}{2})x \, dx.$$

## Example 4 (Robin BCs)

$$-y'' = \lambda y$$
,  $y(0) = 0$ ,  $y(1) + y'(1) = 0$  is an SL problem with:

• Eigenvalues:  $\lambda_n = \omega_n^2$ ,  $n = 1, 2, 3, ..., [\omega_n]$ 's are the positive roots of  $y(x) = x - \tan x$ ].

• Eigenfunctions: 
$$y_n(x) = \sin(\omega_n x)$$
.

The orthogonality of the eigenvectors means that

$$\langle y_m, y_n \rangle := \int_0^1 y_m(x) y_n(x) w(x) \, dx = \int_0^1 \sin(\omega_m x) \, \sin(\omega_n x) \, dx = \begin{cases} 0 & \text{if } m \neq n \\ ??? & \text{if } m = n. \end{cases}$$

Though there isn't a nice closed-form solution, we still have  $||y_n|| := \langle y_n, y_n \rangle^{1/2}$ .

Generalized Fourier series: any function f(x), continuous on [0, 1] satisfying f(0) = 0, f(1) + f'(1) = 0 can be written *uniquely* as

$$f(x) = \sum_{n=1}^{\infty} b_n \sin \omega_n x$$

where

$$b_n = \frac{\langle f, \sin \omega_n x \rangle}{\langle \sin \omega_n x, \sin \omega_n x \rangle} = \frac{\int_0^1 f(x) \sin \omega_n x \, dx}{||\sin \omega_n x||^2} = \frac{\int_0^1 f(x) \sin \omega_n x \, dx}{\int_0^1 (\sin \omega_n x)^2 \, dx}$$